Heavy-Flavor Results from CMS

Keith A. Ulmer*

University of Colorado
on behalf of the CMS Collaboration
E-mail: keith.ulmer@colorado.edu

Heavy-flavor physics offers the opportunity to make indirect tests of physics beyond the Standard Model through precision measurements, and of quantum chromodynamics (QCD) through particle production studies. The rare decays $B_s^0(B^0) \rightarrow \mu^+\mu^-$ and $D^0 \rightarrow \mu^+\mu^-$ are excellent tests of the flavor sector of the Standard Model and are sensitive to new physics. We report on studies of these decays and present the first observation of the excited b baryon Ξ_b^{*0} in strong decays to Ξ_b^- and a charged pion, the observation of two B_c meson decay modes and production properties of the Λ_b baryon, all performed with the CMS experiment in pp collisions at $\sqrt{s} = 7$ TeV.

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*Speaker.

1. Introduction

The CMS experiment at the LHC pursues a diverse program of heavy-flavor measurements. We present here some recent highlights of this program, including searches for the rare decays $B_s(B^0) \to \mu^+\mu^-$ and $D^0 \to \mu^+\mu^-$, which are excellent tests of the flavor sector of the Standard Model with sensitivity to possible physics beyond the Standard Model, and the first observation of the excited b baryon Ξ_b^{*0}. We also present the observation of two decay modes of the B_c meson and production properties of the Λ_b baryon. All measurements were performed with the CMS experiment using pp collision data collected in 2011 at a center-of-mass energy of 7 TeV.

CMS is a general-purpose experiment at the Large Hadron Collider [1]. The inner detector contains a silicon tracker composed of pixel layers at radii less than 15 cm and strip layers out to a radius of 110 cm. The central region $|\eta| < 1$ has 3 layers of pixels and 10 layers of strips. The cylindrical geometry of the central region changes to disks in the $r - z$ plane for the forward region. Each side of the interaction region contains two endcap pixel layers and up to 12 layers of strips. The tracker, PbWO$_4$ electromagnetic calorimeter, and brass-scintillator hadron calorimeter are all immersed in a 3.8 T axial magnetic field. Muons and hadrons are tracked within the pseudorapidity region $|\eta| < 2.4$, with a p_T resolution of about 1.5% for tracks used in the analyses presented here.

2. First observation of Ξ_b^{*0}

The first observation of a new beauty baryon, Ξ_b^{*0}, is reported in decays to $\Xi_b^- \pi^+$ [2]. The precision of the CMS tracker is well utilized in this decay chain, reconstructed with three secondary vertices. First, Λ candidates are formed from oppositely charged proton and pion tracks. The Λ candidates are combined with pion candidates, where the charge of the pion is same as that of the pion in the Λ candidate, to form Ξ^- candidates. The Ξ^- candidates are combined with J/ψ candidates, which are formed from the combination of two oppositely charged muons, to form Ξ_b^- candidates. The muons are also used to trigger the event. Finally, the Ξ_b^- candidates are combined with pion candidates to form Ξ_b^{*0} candidates. A total of thirty variables are used in the final selection, with the most signal-to-background discriminating power coming from displaced-track impact parameter and displaced-vertex requirements. Figure 1 shows the final invariant-mass distributions for the Ξ_b^- and Ξ_b^{*0} candidates. The significance of the signal as measured by the log of the ratio of the fit likelihood modeled with and without a signal is 6.9 standard deviations.

3. Λ_b production

Decays of Λ_b baryons to the final state $J/\psi\Lambda$, with $J/\psi \to \mu^+\mu^-$ and $\Lambda \to p\pi^-$, are used to measure the differential production cross section verses rapidity and transverse momentum [3]. Events are triggered with a displaced dimuon trigger that required dimuon $p_T > 7$ GeV and a separation of the dimuon vertex from the beamspot greater than three times its uncertainty. Oppositely charged muon candidates are combined to form J/ψ candidates, while Λ candidates are formed from oppositely charged pion and proton candidates. The Λ_b candidates are formed by combining a J/ψ candidate with a Λ candidate. The Λ_b yields are obtained by fitting the invariant-mass distribution to a peak above a combinatorial background. The efficiency to reconstruct a Λ_b candidate
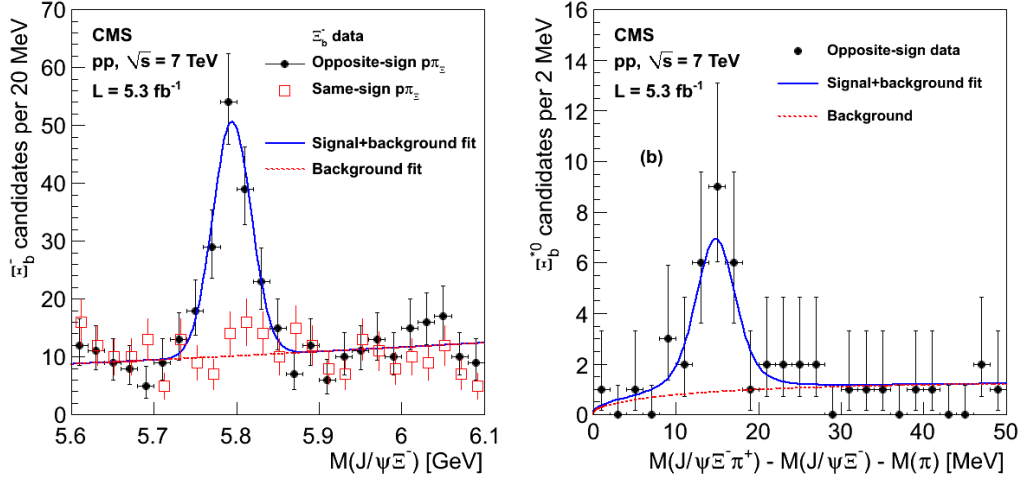
Figure 1: Invariant-mass distributions for Ξ_b^- candidates (left) and Ξ_b^{*0} candidates (right). The dashed (solid) curve represents the background (signal-plus-background) fit function.

is determined from a combination of techniques using Monte Carlo simulation and measurements in the CMS data, with a total efficiency ranging from 0.3% for low-p_T Λ_b candidates to 4.0% for high-p_T candidates. The yields and efficiencies are computed separately in bins of Λ_b rapidity and p_T to obtain a measurement of the differential production cross sections. The production cross section is determined for p_T of the Λ_b starting at 10 GeV and for the central region with the Λ_b rapidity extending up to 2.0. The results are shown in Fig. 2.
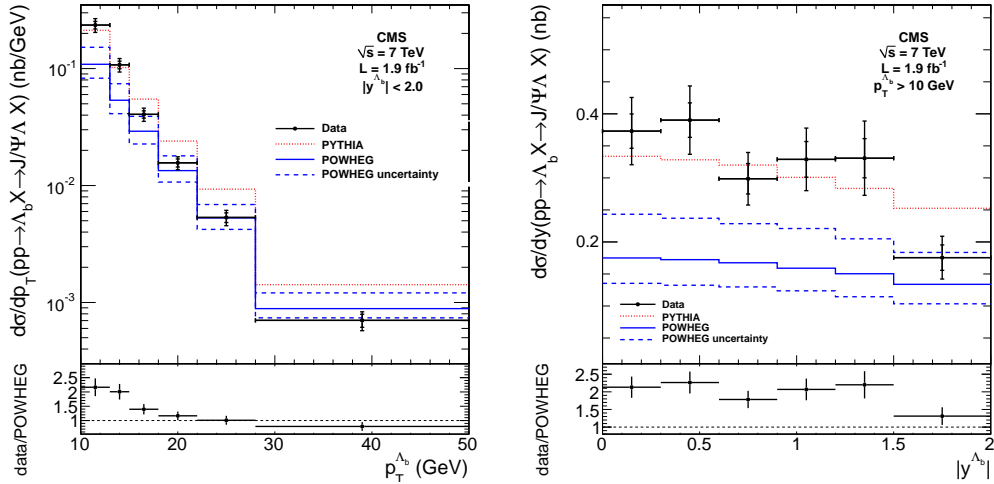

Figure 2: Measured differential cross sections times branching fraction $\mathrm{d}\sigma/\mathrm{d}p_T^{\Lambda_b} \times \mathscr{B}(\Lambda_b \to J/\psi\Lambda)$ (left) and $\mathrm{d}\sigma/\mathrm{d}y^{\Lambda_b} \times \mathscr{B}(\Lambda_b \to J/\psi\Lambda)$ (right) compared to the theoretical predictions from PYTHIA and POWHEG, including theoretical uncertainties on the POWHEG predictions, are shown in the top plots. Ratios of data over the POWHEG predictions are shown in the bottom plots.

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4. Observation of B_c Decay Modes

The B_c^+ meson, consisting of a beauty quark and a charm quark, presents a unique laboratory to study heavy-quark dynamics. We present the observation of two decay channels, $B_c^+ \to J/\psi \pi^+$ and $B_c^+ \to J/\psi \pi^+ \pi^- \pi^+$ [4]. In both channels, J/ψ decays to two muons are used to trigger the events. Either one or three charged pion candidates are combined with the J/ψ candidate to form B_c^+ candidates. A total yield of 330 ± 36 events are observed for the one-pion-decay channel, while 108 ± 19 events are observed for the three-pion-decay channel. Figure 3 (left) shows the invariant-mass distribution for the three-pion-decay channel.
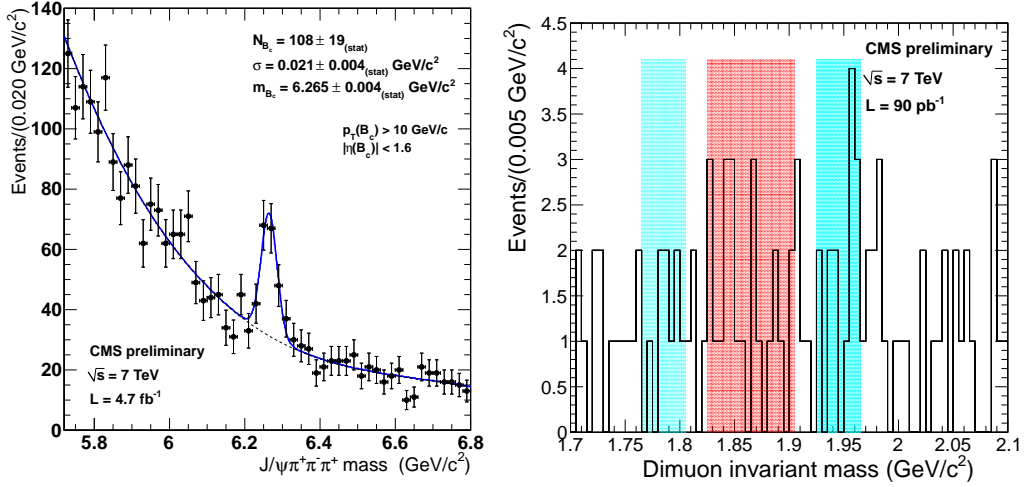

Figure 3: Invariant-mass distribution for B_c candidates in the $J/\psi \pi^+ \pi^- \pi^+$ channel (left), where the dashed (solid) line represents the background (signal-plus-background) fit function, and $D^0 \to \mu^+ \mu^-$ candidates (right), where the signal (control) region is shown shaded in red (blue).

5. Search for $D^0 \to \mu^+ \mu^-$

The rare flavor-changing neutral-current decay $D^0 \to \mu^+ \mu^-$ is highly suppressed in the Standard Model, but can be enhanced in new-physics scenarios. We present results from the first search for this decay at CMS based on a small portion of the total data set representing $\approx$90 pb^{-1} [5]. A single-muon trigger is used to search for the signal decay and for a control channel, $D^0 \to K^- \mu^+ \nu$, which has similar kinematics. Many systematic effects cancel in the ratio of the efficiencies for the signal and control channels. Signal candidates are reconstructed by combining two oppositely charged muon candidates. The dimuon candidate momentum is required to point back to the beamline. The control channel is reconstructed by identifying D^0 candidates from $D^{*+} \to D^0 \pi^+$ decays. Backgrounds for the search sample are measured by interpolating from the dimuon mass sidebands in data on either side of the D^0 mass range. No evidence of $D^0 \to \mu^+ \mu^-$ decays is observed, and the control sample yield and known branching fraction are used to determine an upper limit on the branching fraction. At 90% confidence level, an upper limit on the $D^0 \to \mu^+ \mu^-$ branching fraction

is calculated with the CL_s technique as $< 5.4 \times 10^{-7}$. Figure 3 (right) shows the invariant-mass distribution for the $D^0 \to \mu^+ \mu^-$ candidates, where the red and blue regions are the signal and control regions, respectively.

6. Search for $B^0_{(s)} \to \mu^+ \mu^-$

The rare flavor-changing neutral-current decays $B^0_{(s)} \to \mu^+ \mu^-$ are highly suppressed in the Standard Model, but can be altered in many new-physics scenarios. In particular, limits at the LHC are now approaching the Standard Model prediction for $B^0_s \to \mu^+ \mu^-$ of $(3.2 \pm 0.2) \times 10^{-9}$ [6]. We present the results of a simultaneous search for $B^0 \to \mu^+ \mu^-$ and $B^0_s \to \mu^+ \mu^-$ at CMS in pp collisions at $\sqrt{s} = 7 \, \text{TeV}$, with a data sample corresponding to an integrated luminosity of $5 \, \text{fb}^{-1}$ [7]. Candidate $B^0_{(s)}$ mesons are reconstructed by fitting oppositely charged muon candidates to a common vertex. Signal events are selected preferentially by requiring the dimuon vertex to be well separated from the beamspot, the dimuon candidate to be well isolated from other detector activity and that the two muons form a good vertex.

The search is performed by counting events with dimuon invariant mass within appropriate windows for $B^0_{(s)}$ mesons. The event sample is divided into two categories, barrel and endcap, where the former contains only candidates when both muons are within $|\eta| < 1.4$, which have the best momentum resolution. Backgrounds arise from combinatorial muon pairs, which are measured by extrapolating from the sidebands of the dimuon invariant-mass distributions, and from peaking backgrounds from semileptonic or hadronic B decays where hadrons are misidentified as muons. The peaking background is measured by scaling the background expectation from simulation by the hadron misidentification rates measured in CMS data. No significant signal is observed in either channel. The dimuon invariant-mass distributions are shown in the top plots of Fig. 4.

Upper limits on the branching fractions are determined through a normalization channel, $B^+ \to J/\psi K^+$. Many systematic uncertainties cancel in the ratio of the yields and efficiencies between the signal and normalization channels. Upper limits of $\mathscr{B}(B^0_s \to \mu^+ \mu^-) < 7.7 \times 10^{-9}$ and $\mathscr{B}(B^0 \to \mu^+ \mu^-) < 1.8 \times 10^{-9}$ are set at 95% confidence level using the CL_s technique. The $\text{CL}_\text{s+b}$ curve for $B^0_s \to \mu^+ \mu^-$ is shown in the lower plot of Fig. 4. Additionally, a combination of the CMS result with measurements of the same channels from LHCb and ATLAS is performed, which provides the best upper limits to date for these channels, $\mathscr{B}(B^0_s \to \mu^+ \mu^-) < 4.2 \times 10^{-9}$ and $\mathscr{B}(B^0 \to \mu^+ \mu^-) < 8.1 \times 10^{-10}$ at 95% confidence level [8].

7. Conclusion

A selection of results from the CMS heavy-flavor physics program has been shown, including new limits on searches for rare decays and new measurements of heavy-flavor production.

References

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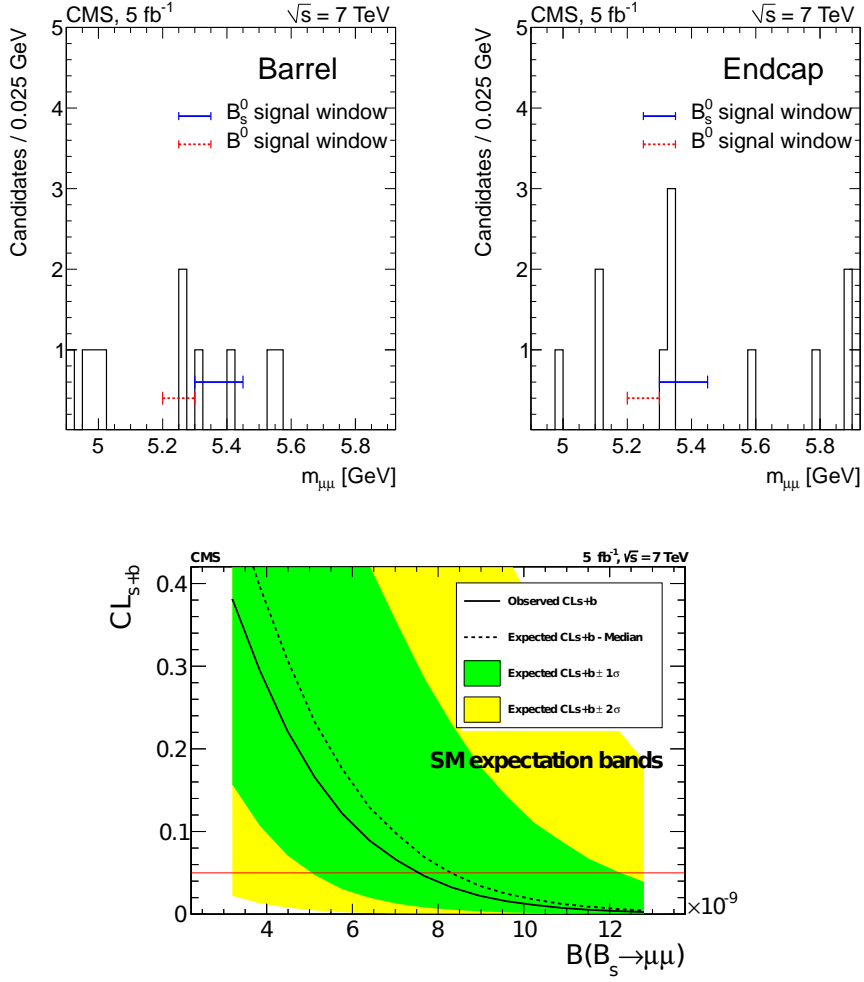[3] S. Chatrchyan *et al.* [CMS Collaboration], Phys. Lett. B **714**, 136 (2012).

Figure 4: Invariant-mass distributions for $B^0_{(s)} \to \mu^+\mu^-$ candidates in the barrel (upper left) and endcap (upper right) regions. The blue and red regions are the signal windows for B^0_s and B^0, respectively. The observed (solid line) and expected (dotted line) CL_{s+b} curves for $B^0_s \to \mu^+\mu^-$ compared to SM expectation (bottom). The green and yellow regions show the ± 1 and ± 2 standard deviation variations in the expected sensitivities.

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